The initial conditions are

$$u(r,o) = \frac{\partial u(r,o)}{\partial r} = o$$
(14)

Our approach in the following derivations is first to obtain a solution for the case of step application of microwave energy at some instant t = o and then to extend the solution to a square pulse using Duhamel's theorem [16].

## Step Function Excitation

If we write the displacement u(r,t) as

$$u(r,t) = u_{e}(r) + u_{t}(r,t)$$
 (15)

and substitute equation (15) into equation (8), the equation of motion becomes two differential equations: a stationary one and a time-varying one. Thus,

$$\frac{d^2 u_s(r)}{dr^2} + \frac{2}{r} \frac{d u_s(r)}{dr} - \frac{2}{r^2} u_s(r) = u_0 F_r(r)$$
(16)

$$\frac{\partial^2 u_t(r,t)}{\partial r^2} + \frac{2}{r} \frac{\partial u_t(r,t)}{\partial r} - \frac{2}{r^2} u_t(r,t) = \frac{1}{c_1^2} \frac{\partial^2 u_t(r,t)}{\partial t^2}$$
(17)

The corresponding boundary conditions at r = a are

$$(\lambda + 2\mu) \frac{du}{dr} + 2\lambda \frac{u}{r} = 0$$
(18)

and

$$(\lambda + 2\mu) \frac{\partial u_t}{\partial r} + 2\lambda \frac{u_t}{r} = 0$$
(19)

The solution of equation (16), using the boundary condition of equation (18), is given by

$$u_{s}(r) = u_{0}\left[\frac{a}{N\pi} j_{1}\left(\frac{N\pi r}{a}\right) + \frac{4\mu}{3\lambda + 2\mu} \frac{\gamma}{N^{2}\pi^{2}}\right], \quad N = \begin{cases} 1,3,5...\\ 0,2,4... \end{cases}$$
(20)

where  $j_1(\frac{N\pi r}{a})$  is the spherical Bessel function.

Now we let

$$u_{t}(r,t) = R(r) T(t)$$
 (21)

and use the method of separation of variables to solve equation (17) for the time-varying component. Inserting equation (21) into equation (17) yields the two ordinary differential equations:

$$\frac{d^2 R}{dr^2} + \frac{2}{r} \frac{dR}{dr} + (k^2 - \frac{2}{r^2})R = 0$$
(22)